# Traveling fronts for lattice neural field equations

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### Excitatory neural network on a lattice

$$\dot{u}_n(t) = -u_n(t) + \sum_{j \in \mathbb{Z}} K_j S(u_{n-j}(t)), \quad (n,t) \in \mathbb{Z} \times (0,\infty)$$

- ▶  $u_n(t) \in \mathbb{R}$ : membrane potential of neuron labelled *n* at time *t*;
- $K_j \ge 0$ : strength of interactions (all to all infinite range)
- $u \mapsto S(u)$ : firing rate function (of sigmoidal type)
- discrete version of

$$\partial_t \mathbf{u}(x,t) = -\mathbf{u}(x,t) + \int_{\mathbb{R}} \mathcal{K}(y) S(\mathbf{u}(x-y,t)) dy$$

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References: Wilson-Cowan '72, Amari '77, Hopfield'84, Ermentrout '98, Coombes '05, Bressloff '12

### Traveling waves

Traveling wave: speed  $c \in \mathbb{R}$ 

$$u_n(t) = \mathbf{u}(n-ct), \quad \mathbf{u}: \mathbb{R} \to \mathbb{R}$$

Traveling wave equation: x = n - ct

$$-c\mathbf{u}'(x) = -\mathbf{u}(x) + \sum_{j \in \mathbb{Z}} K_j S(\mathbf{u}(x-j)), \quad x \in \mathbb{R},$$
(1a)  
$$\lim_{x \to -\infty} \mathbf{u}(x) = 1 \text{ and } \lim_{x \to +\infty} \mathbf{u}(x) = 0,$$
(1b)

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- FDE of mixed type (ill posed as a Cauchy problem)
- dynamics depend on (all) past  $j \leq 0$  and (all) future  $0 \leq j$

It is assumed that when c = 0 equation (1) is an infinite recurrence.

#### Assumptions

On the firing rate function S:

**Hypothesis (H1)** - **Bistable nonlinearity.** We suppose that:

(i)  $S \in \mathscr{C}'_b(\mathbb{R})$  for  $r \ge 2$  with S(0) = 0 and S(1) = 1 together with S'(0) < 1and S'(1) < 1;

- (ii) there exists a unique  $\theta \in (0, 1)$  such that  $S(\theta) = \theta$  with  $S'(\theta) > 1$ ;
- (iii)  $u \mapsto S(u)$  is strictly nondecreasing on [0, 1] and there exists  $s_m > 1 > s_0 > 0$  such that  $s_0 < S'(u) \le s_m$  for all  $u \in [0, 1]$ .

On the weights  $(K_j)_{j \in \mathbb{Z}}$ :

Hypothesis (H2) - Weights. We suppose that:

(i) 
$$\sum_{n \in \mathbb{Z}} K_n = 1$$
;  
(ii) for all  $n \in \mathbb{Z}$ , we have  $K_n = K_{-n} \ge 0$  and  $K_{\pm 1} > 0$   
(iii)  $\sum_{n \in \mathbb{Z}} |n| K_n < \infty$ .

# Existence of monotone traveling waves

#### Theorem

Suppose that the Hypotheses (H1)-(H2) are satisfied then there exists a traveling wave solution  $u_n(t) = \mathbf{u}_*(n - c_*t)$  of LNFE such that the profile  $\mathbf{u}_*$  satisfies the traveling wave problem (1). Moreover:

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a strictly decreasing sequence.

### Sketch of the proof

The key idea (Bates & Chen '99) is to regularize the traveling wave equation

$$-c\mathbf{u}'(x) = -\mathbf{u}(x) + \sum_{j \in \mathbb{Z}} K_j S(\mathbf{u}(x-j))$$

Let  $\Psi\in \mathscr{C}^\infty(\mathbb{R}), \ \Psi\geq 0, \ \int_{\mathbb{R}}\Psi(x)\mathrm{d}x=1$ , even and with compact support

$$\rho_m(x) := m\Psi(mx), \quad \mathcal{K}_m(x) := \sum_{j=-m}^m \frac{1}{\omega_m} K_j \rho_m(x-j), \text{ with } \omega_m := \sum_{j=-m}^m K_j$$

New traveling wave problem

$$-c_m \mathbf{u}'_m = -\mathbf{u}_m + \mathcal{K}_m * S(\mathbf{u}_m), \text{ on } \mathbb{R},$$
(2a)

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$$\lim_{x \to -\infty} \mathbf{u}_m(x) = 1 \text{ and } \lim_{x \to +\infty} \mathbf{u}_m(x) = 0.$$
 (2b)

Apply the results of Ermentrout & McLeod '93 to get a monotone solution  $(\mathbf{u}_m, c_m)$  of (2) then pass to the limit  $m \to +\infty$ .

# Uniqueness of traveling waves with nonzero speed

#### Theorem

Let  $(\mathbf{u}_*, c_*)$  be a solution to (1) as given in Theorem 1, such that  $c_* \neq 0$ . Let  $(\hat{\mathbf{u}}, \hat{\mathbf{c}})$  be another solution to (1). Then  $c = \hat{c}$  and, up to a translation,  $\mathbf{u}_* = \hat{\mathbf{u}}$ .

Ideas of the proof:

Construct appropriate sub and super solutions of the form

$$w_n^{\pm}(t) := \mathbf{u}_* \left( n - c_* t + \xi_0 \mp \sigma \gamma (1 - e^{-\beta t}) \right) \pm \gamma e^{-\beta t}, \quad \forall n \in \mathbb{Z}$$

for some parameters  $\xi_0$ ,  $\sigma$ ,  $\gamma$  and  $2\beta = \min \{1 - S'(0); 1 - S'(1)\}$ .

Use comparison principle and a "squeezing" technique to prove uniqueness

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Reference: Chen'97

# Spectral stability

LNFE in moving coordinate:

$$-c\mathbf{u}'(x) = -\mathbf{u}(x) + \sum_{j \in \mathbb{Z}} K_j S(\mathbf{u}(x-j))$$

Linearized operator around  $(\mathbf{u}_*, c_*)$  with  $c_* \neq 0$ :

$$\mathcal{L}: H^1(\mathbb{R}) o L^2(\mathbb{R}), \qquad \mathcal{L} \mathbf{v} := c_* \mathbf{v}' - \mathbf{v} + \mathcal{K}_\delta * [S'(\mathbf{u}_*) \mathbf{v}]$$

where  $\mathcal{K}_{\delta} * \mathbf{w} = \sum_{j \in \mathbb{Z}} \mathcal{K}_{j} \mathbf{w}(\cdot - j).$ 

Floquet-like spectral structure:

$$\mathcal{L}(e^{2\pi i x} \mathbf{u}) = e^{2\pi i x} (2\pi i c_* + \mathcal{L}) \mathbf{u}$$

- spectrum invariant under shifts by  $2\pi i c_*$
- lattice doesn't feel oscillations on scale smaller than distance in lattice

# Spectral properties of $\mathcal{L}$

#### **Hypothesis** (H2 $\eta$ ) - **Exponential localization**. We suppose that:

- (i)  $(K_j)_{j\in\mathbb{Z}}$  satisfies (H2);
- (ii) there exists  $\eta > 0$ , such that  $\sum_{j \in \mathbb{Z}} K_j e^{\eta |j|} < \infty$ .

#### Proposition

Assume that Hypotheses (H1)-(H2 $\eta$ ) and that ( $\mathbf{u}_*, c_*$ ) is the traveling wave solution given in Theorem 1 with  $c_*$ . We have:

- $\blacktriangleright$  0 is an algebraically simple eigenvalue of  ${\cal L}$  with a negative eigenfunction  $u'_*;$
- ► the adjoint operator L<sup>\*</sup> has a negative eigenfunction, denoted q ∈ C<sup>1</sup>(R), corresponding to the simple eigenvalue 0;
- for all  $0 < \kappa < 2\beta$  the operator  $\mathcal{L} \lambda$  is invertible as an operator from  $H^1(\mathbb{R})$  to  $L^2(\mathbb{R})$  for all  $\lambda \in \mathbb{C} \setminus 2\pi i c_* \mathbb{Z}$  such that  $\Re(\lambda) \ge -\kappa$ ;
- ▶ there exist  $\eta_*, \eta_{**} \in (0, \eta)$  and some constants  $C_* > 0$ ,  $C_{**} > 0$  such that

$$|\mathbf{u}_{*}'(x)| \leq C_{*}e^{-\eta_{*}|x|}\|\mathbf{u}_{*}\|_{L^{\infty}(\mathbb{R})}, \text{ and } |\mathbf{q}(x)| \leq C_{**}e^{-\eta_{**}|x|}\|\mathbf{q}\|_{L^{\infty}(\mathbb{R})}.$$

# Toward nonlinear stability

GOAL: given existence and spectral stability, prove nonlinear stability

LNFE can be written

$$\dot{\mathfrak{u}}(t) = \mathcal{F}(\mathfrak{u}(t)), \text{ with } \mathfrak{u} = (u_n)_{n \in \mathbb{Z}} \in \mathbb{R}^{\mathbb{Z}}$$

where  $\bar{\mathfrak{u}}(t) = (\mathfrak{u}_*(n-c_*t))_{n\in\mathbb{Z}}$  with  $c_* \neq 0$  is a solution with

$$\bar{\mathfrak{u}}_n(t) = \bar{\mathfrak{u}}_{n-1}\left(t-\frac{1}{c_*}\right), \quad n \in \mathbb{Z}$$

- TW is relative periodic orbit
- Linearization:

$$\dot{\mathfrak{v}}(t) = D\mathcal{F}(\bar{\mathfrak{u}}(t))\mathfrak{v}(t)$$

 $\Rightarrow$  use spectral information on  $\mathcal{L}$  to obtain decay estimates on the flow of  $\dot{\mathfrak{v}}(t) = D\mathcal{F}(\bar{\mathfrak{u}}(t))\mathfrak{v}(t)$  and prove nonlinear stability

### Main strategy

1. Moving coordinate frame:

$$\partial_t \mathbf{v} = \mathcal{L} \mathbf{v}, \quad (\mathcal{L} - \lambda) \mathbf{G} = \delta(\cdot - x_0),$$

- spectral properties on  $\mathcal{L}$
- resolvent kernel  $\mathbf{G}_{\lambda}(x, x_0)$
- 2. Coordinates of the lattice:

$$\dot{\mathfrak{v}}(t)=D\mathcal{F}(ar{\mathfrak{u}}(t))\mathfrak{v}(t), \hspace{1em} \mathfrak{v}(t_0)=(\delta(j-j_0))_{j\in\mathbb{Z}} \hspace{1em}\Rightarrow\hspace{1em} ext{solution}\hspace{1em} \mathfrak{v}^{t_0j_0}(t)$$

Green's function

$$\mathcal{G}_{jj_0}(t,t_0) = \mathfrak{v}_j^{t_0j_0}(t)$$

Relationship (Benzoni-Gavage, Huot, Rousset '03)

$$\mathcal{G}_{jj_0}(t,t_0) = \frac{1}{2\pi \mathbf{i}} \int_{R-\mathbf{i}\pi c_*}^{R+\mathbf{i}\pi c_*} e^{\lambda(t-t_0)} \mathbf{G}_{\lambda}(j-c_*t,j_0-c_*t_0) \mathrm{d}\lambda \quad R \gg 1$$

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# Green's function representation

#### Proposition

Assume that Hypotheses (H1)-(H2 $\eta$ ) and that ( $\mathbf{u}_*, c_*$ ) is the traveling wave solution given in Theorem 1 with  $c_*$ . Then there exists  $\epsilon > 0$  such that for all  $\lambda \in \mathbb{C}$  with  $0 < |\lambda| < \epsilon$ , we have the representation

$$\mathbf{G}_{\lambda}(x, x_0) = \mathbf{E}_{\lambda}(x, x_0) + \widetilde{\mathbf{G}}_{\lambda}(x, x_0),$$

where  $\mathbf{E}_{\lambda}$  can be written as

$$\mathsf{E}_{\lambda}(x,x_{0}) = -\frac{1}{\lambda \int_{\mathbb{R}} \mathsf{q}(z) \mathsf{u}_{*}'(z) \mathrm{d}z} \mathsf{u}_{*}'(x) \mathsf{q}(x_{0}),$$

while the remainder term depends analytically of  $\lambda$  in the region  $|\lambda| < \epsilon$ .

We need estimates on  $\widetilde{\boldsymbol{\mathsf{G}}}_{\lambda}$  and we expect that

$$\left|\widetilde{\mathbf{G}}_{\lambda}(x,x_{0})\right| \leq Ce^{-\omega'|x-x_{0}|}, \quad \text{(ongoing work)}$$

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for some C > 0 and  $\omega' > 0$ .

Related works: Beck etal '10, Hupkes-Sandstede '13, Schouten-Hupkes '17

### Discussion

For Lattice Neural Field Equation

$$\dot{u}_n(t) = -u_n(t) + \sum_{j\in\mathbb{Z}} \mathcal{K}_j S(u_{n-j}(t))$$

- existence & uniqueness (up to translation) of monotone traveling front solutions for "bistable" type of kinetics
- spectral stability of nonzero wave speed traveling fronts (need an exponential localization of interactions)
- toward a nonlinear stability result study of pointwise Green's functions
- difficulty & novelty: infinite nonlinear range interactions
- study numerical approximation schemes of continuous NFE
- study other type of networks and/or kinetics (monostable, with linear adaptation)



# Thematic semester on Wathematics Computer science and biology

**Deterministic and Stochastic Models** in Neurosciences

